

Simulation of heat flux through multilayer structures

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R. MacKenzie gratefully acknowledges the support of the **Engineering and Physical Sciences Research Council** (EPSRC), U.K.



We gratefully acknowledge the EC IST project **FAST ACCESS** (IST-004772).



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Presentation outline



- Heat flow through structures with multiple epitaxial layers
 - > Thermal conductivity of structures with multiple interfaces
 - > Thermal boundary resistance (TBR)
- \cdot Heat flow in 1.3 μ m dilute nitride edge-emitting lasers
 - Full electro-opto-thermal simulations
 - Introduction of TBR into the model
- \cdot Discussion of effect in other (dilute nitride) devices
- · Conclusions

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Thermal conductivity of superlattices



- GaAs/AlAs superlattices have a much lower thermal conductivity than one would predict from the bulk values alone.¹ (3x-10x lower)
 - Bulk GaAs/AlAs thermal conductivity = 58.4m⁻¹K⁻¹
 - Superlattice thermal conductivity = 5.0m⁻¹K⁻¹



- This effect is mainly due to phonon scattering/reflections at material interfaces
- TBR first observed by Kapitza (1941)²



- i = incident wave
- r = reflected wave
- t = transmitted wave

[1] W.S. Capinski *et. al.*, Phys. Rev. B Vol. 59, No. 12, p.8105 (1999).
[2] Collected papers of P.L. Kapitza, Vol. 2, Pergamon, Oxford, p. 581 (1965).

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How does structure size affect the conductivity?



Consider a superlattice with a period L, where Λ is the average phonon mean free path (20-140nm)

One can distinguish two regimes:

- 1) $L \approx \Lambda$ A bulk thermal conductivity can be used between the interfaces by placing a thermal resistance at each boundary (TBR)
- 2) $L \ll \Lambda$ The situation becomes more complicated with phonons reflecting off multiple layers and gaps forming in the dispersion relations

> Edge-emitting lasers fall within the L $\approx \Lambda$ regime

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How does TBR affect dilute nitride EELs?





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What values of TBR should be used?



- Values of TBR are depend on:
 - > The acoustic mismatch of the materials
 - Masses Elastic constants -> Speed of sound in materials
 - Similar to Snell's law
 - > The quality of epitaxial interfaces
 - E.g. nitrogen plasma damage (dilute nitride devices)
 - Layer thickness
- Exhaustive experimental characterization of the effect is not complete
 Still no real consensus on microscopic models for TBR
- Diffuse mismatch model is used in this work
 - Has shown some agreement with experiment
 - A range of values will be used to examine the impact of TBR



Interface	TBR value (m²K/W)	Method	Device
GaAs/AlGaAs/InGaAsN	≈1.2x10 ⁻⁹	DMM	EEL
GaN/Si¹	7x10 ⁻⁸	experiment	HEMT
GaN/SiC ¹	1.2x10 ⁻⁷	experiment	HEMT
AIN/Si ¹	7-8x10 ⁻⁸	experiment	Thin film

- TBR is a known problem in nitride HEMTs
- TBR is less well studied in EELs

1) J. Kuzmík et.al., J. Appl. Phys. Vol. 101, 054508 (2007).

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Device simulator



Electro-thermal Model Bipolar 2D Drift Diffusion (DD) model (0th and 1st moments of the Boltzmann Transport Equation (BTE)) Poisson's equation QW capture/escape equations for each QW 2D lattice heat equation Heat sources derived from 2nd moment of BTE

Optical Model

- Photon rate equation
- Valance band structure calculated using 4x4 band k.p
- Band anti-crossing model for the conduction band
- Fermi's Golden rule used to calculate stimulated/spontaneous emission rates

All equations solved using Newton's method



Discretization scheme for inclusion of TBR



• The lattice heat equation is commonly solved in thermal models:

$$\rho_L C_L \frac{\partial T}{\partial t} = \nabla \cdot \left(k \nabla T \right) + H$$

 However, because of abrupt thermal resistances at epitaxial interfaces one must solve:

(1)
$$\left(\frac{\partial T}{\partial x}\right)_{r+1/2}^{3}k_{1} = k_{2}\left(\frac{\partial T}{\partial x}\right)_{r+1/2}^{4}$$

 Introduce a step in temperature proportional to the boundary resistance:

(2)
$$T_{r+1/2}^{(3)} - T_{r+1/2}^{(4)} = Rk_2 \left(\frac{\partial T}{\partial x}\right)_{r+1/2}^3$$

 Adapted from a scheme to model discontinuities Quasi-TE modes of semiconductor waveguides¹

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Thermal profile with and with out TBR





- Up to half a degree difference in peak temperature of device
- Small temperature differences are important for accurate models
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15KV X20,000 6mm



Difference between simulations with and without TBR



- Joule heating and free carrier absorption in ridge
 - High heat flux out of ridge
 - > TBR has a large effect at bottom of ridge

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Change in QW temperature due to TBR vs. injection current



• The interfaces introduce a small increase in QW temperature

Impact of TBR on QW temperature



Temperature of QW for various TBR values

Difference in QW temperature due to TBR



- A small increase in QW temperature can be seen
 - Although small, the effect may need to be included in some cases

Impact of TBR on output power



L-I curves for isothermal, ordinary thermal and thermal with TBR

Difference in L-I curves compared to ordinary thermal case



- x1,x2,x4 and x8 times the value predicted by TBR
- A small decrease in power due to the temperature increase can be seen

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Conclusions



- TBR has been shown to increase the predicted temperature of a $1.3 \mu m$ dilute nitride EEL by up to 0.5K
- A small decrease in optical power is also predicted
- Impact of TBR increased by
 - More interfaces
 - Materials with large acoustic mismatch (GaN/SiC HFETs)
 - Interface defects (nitrogen plasma damage)
 - Increased heat flux
- Need for more *more accurate* TBR values
 - > Ideally from experiment
 - Better numerical models for calculation of TBR are also needed
- For the first time the impact of TBR has been considered within a full electro-optical-thermal laser simulation tool

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Other devices structures where TBR has an impact



- Structures with relatively thick layers (L $\approx \Lambda$)
 - Long wavelength VCSELs mirrors
 - e.g. 1.3µm structures for dilute nitride VCSELs
 - > 30-60 periods
 - Possible increase of up to 5K
 - Carrier heat flux reduces impact of TBR

· Devices with large acoustic mismatch

- GaN/sapphire, GaN/SiC and GaN/AIN
 - > Heterostructure field-effect transistors (HFETs)
- Structures with layer thicknesses that are much smaller than the phonon m.p.f. (L $\ll \Lambda$)
 - Short wavelength VCSELs
 - Quantum cascade lasers
 - > High temperature sensitivity